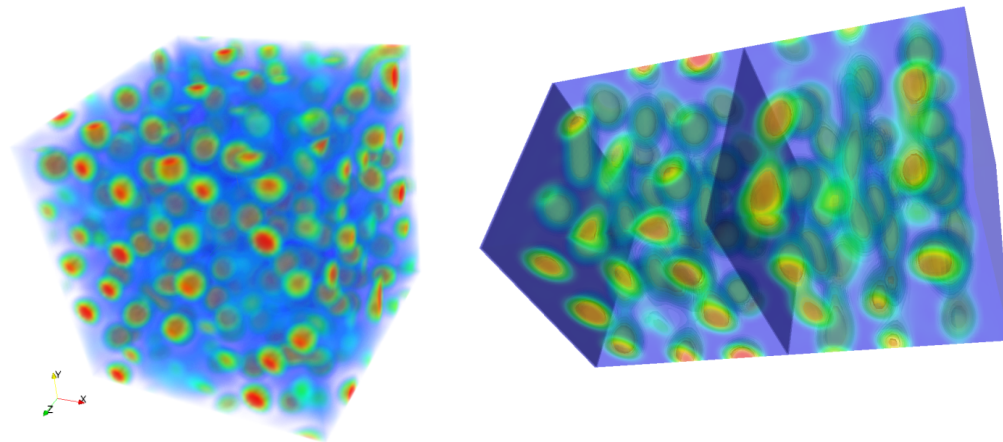


A Finite Element-Based Phase Field Model

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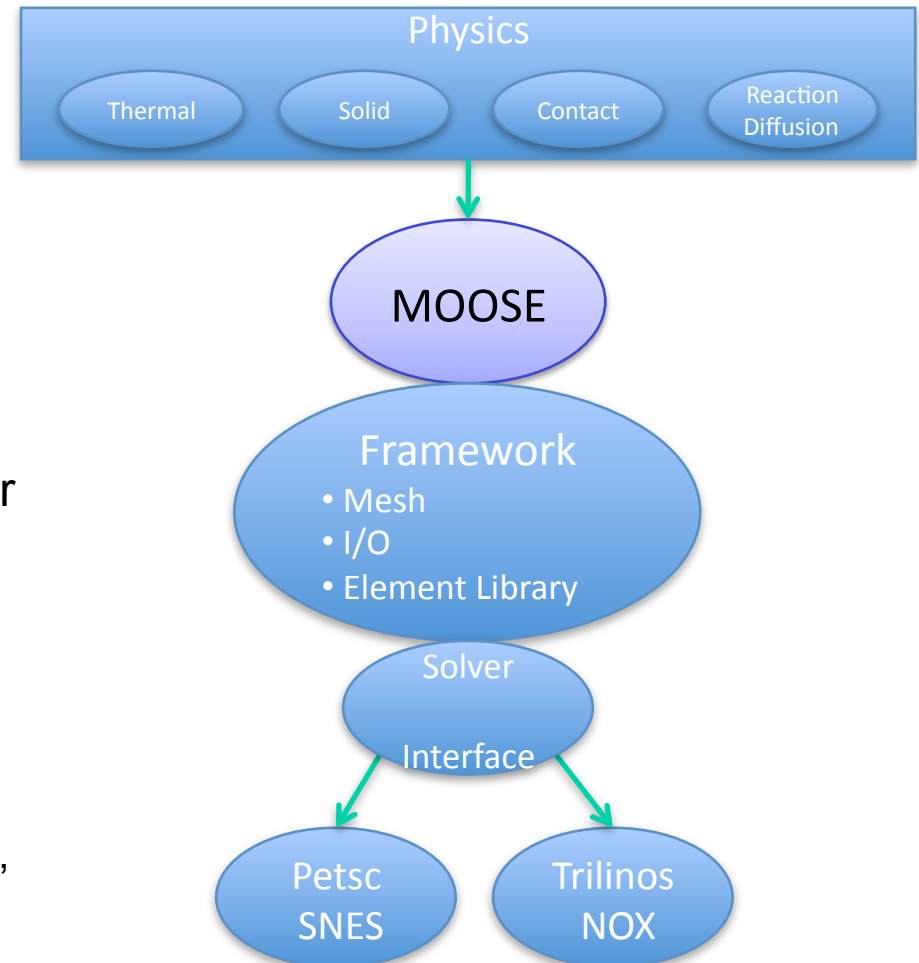


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Multiscale Object Oriented Simulation Environment (MOOSE)

- Plug-and-play modules
 - Simplified coupling
- MOOSE provides a set of interfaces
- Framework provides core set of common services
 - libMesh: <http://libmesh.sf.net>
- Solver Interface abstracts specific solver implementations.
 - Common interface to linear and non-linear solvers
 - More flexible
- Utilize state-of-the-art linear and non-linear solvers
 - Robust solvers are key for “ease of use”



Finite Element Solution of Phase Field Equations

- Strong Form

$$\mathbf{R}_{c_i} = \frac{\partial c_i}{\partial t} - \nabla \cdot \left(M_{ij} \nabla \left(\frac{\partial g_0}{\partial c_i} - \kappa \nabla^2 c_i + \frac{\partial E_{el}}{\partial c_i} \right) \right) = 0$$

$$\mathbf{R}_{\eta_i} = \frac{\partial \eta_i}{\partial t} - L_i \left(\frac{\partial f_0}{\partial \eta_i} - \kappa \nabla^2 \eta_i + \frac{\partial E_{el}}{\partial \eta_i} \right) = 0$$

$$\mathbf{R}_{\mathbf{u}} = \nabla \cdot (\mathbf{C} \nabla \mathbf{u}) - \nabla \cdot (\mathbf{C} \boldsymbol{\epsilon}^*) = 0$$

$$\mathbf{R}_T = \nabla \cdot (k \nabla T) = 0$$
- Weak Form

$$\mathbf{R}_{c_i} = \left(\frac{\partial c_i}{\partial t}, \phi_i \right) + \left(M_{ij} \nabla \frac{\partial g_0}{\partial c_i}, \nabla \phi_i \right) + \left(\kappa \nabla^2 c_i, \nabla \cdot (M_{ij} \nabla \phi_i) \right) + \left(M_{ij} \nabla \frac{\partial E_{el}}{\partial c_i}, \nabla \phi_i \right) = 0$$

$$\mathbf{R}_{\eta_i} = \left(\frac{\partial \eta_i}{\partial t}, \phi_i \right) + L_i \left(\frac{\partial f_0}{\partial \eta_i}, \phi_i \right) + \kappa (\nabla \eta_i, \nabla \phi_i) + \left(\frac{\partial E_{el}}{\partial \eta_i}, \phi_i \right) = 0$$

$$\mathbf{R}_{\mathbf{u}} = (\mathbf{C} \nabla \mathbf{u}, \nabla \phi_i) - (\nabla \cdot \mathbf{C} \boldsymbol{\epsilon}^*, \phi_i) = 0$$

$$\mathbf{R}_T = (k \nabla T, \nabla \phi_i) = 0$$

- FEM discretization

$$c_i(\mathbf{r}) = \sum_{j=1}^N c_i^j \phi_j(\mathbf{r})$$

Discretized using 3rd
order Hermite element

2D: 20 DOF

3D: 36 DOF

$$\eta_i(\mathbf{r}) = \sum_{j=1}^N \eta_i^j \phi_j(\mathbf{r})$$

Discretized using 1st order
Lagrange elements

2D: 8 DOF

3D: 12 DOF

$$\mathbf{u}(\mathbf{r}) = \sum_{j=1}^N \mathbf{u}^j \phi_j(\mathbf{r})$$

Discretized using 1st order
Lagrange elements

2D: 8 DOF

3D: 12 DOF

$$T(\mathbf{r}) = \sum_{j=1}^N T^j \phi_j(\mathbf{r})$$

Discretized using 1st order
Lagrange elements

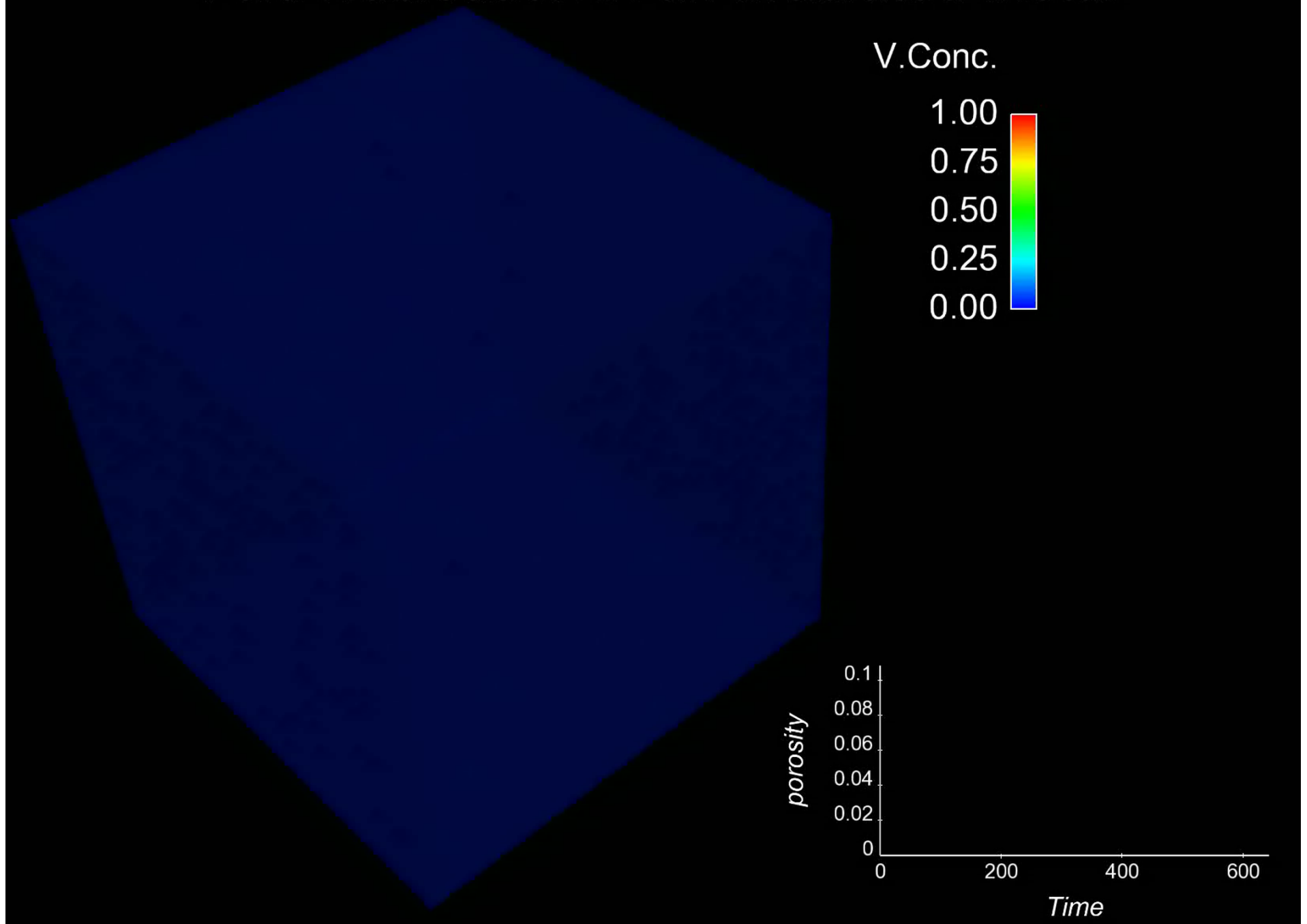
2D: 8 DOF

3D: 12 DOF

MARMOT Summary

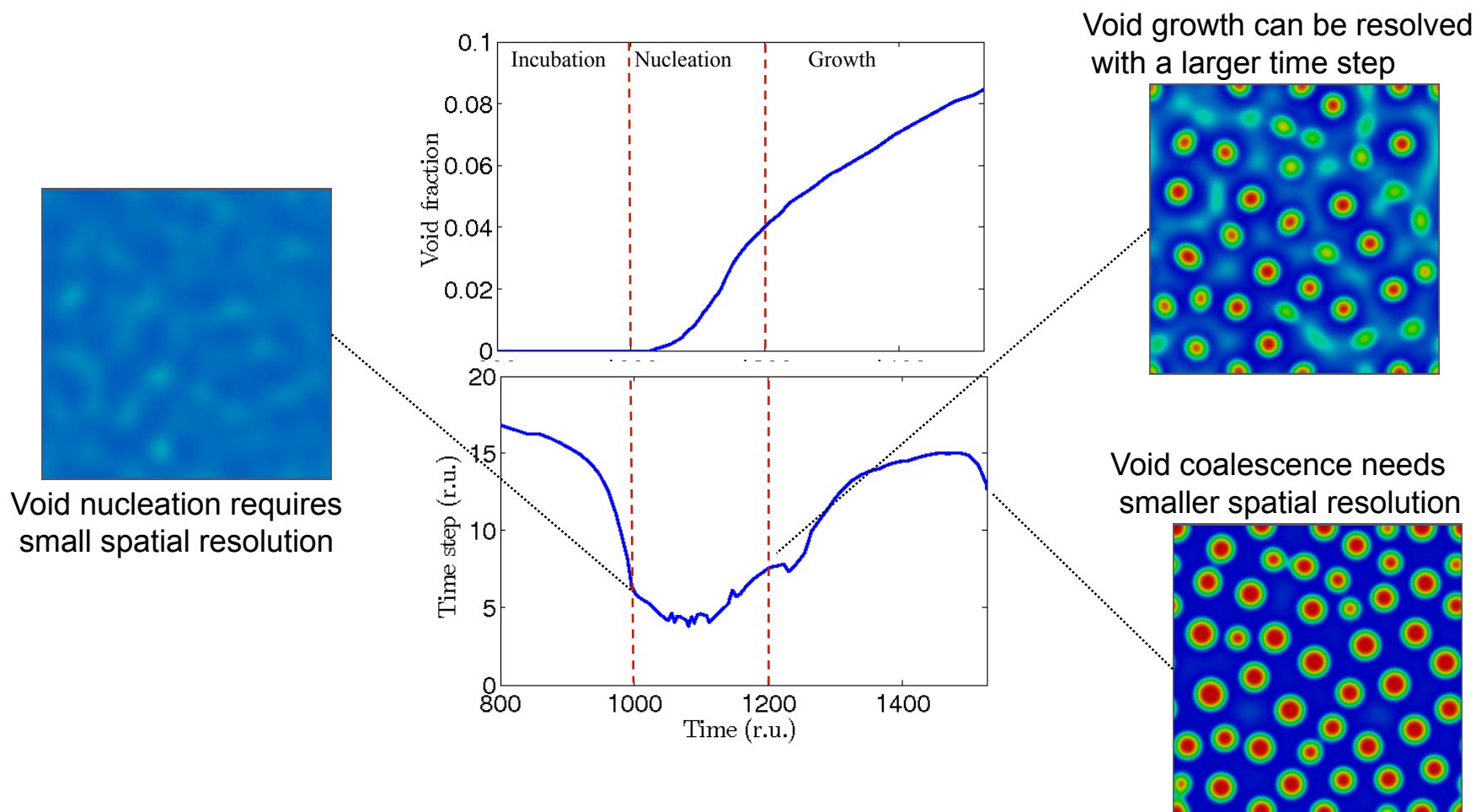
- Implicit solution of the phase field equations using preconditioned JFNK
 - Void nucleation and growth in an irradiated single component metal
 - Grain growth and GB sink effects
- Fully-coupled multiphysics – Heat conduction and linear elastic solid mechanics
- Calculation of bulk properties such as effective thermal conductivity and porosity
- Mesh adaptivity
- Time step adaptivity

Void Nucleation in an Irradiated Metal

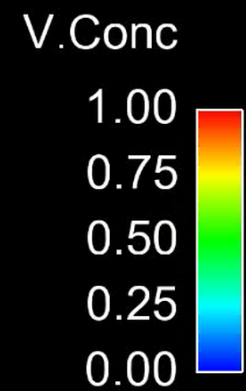
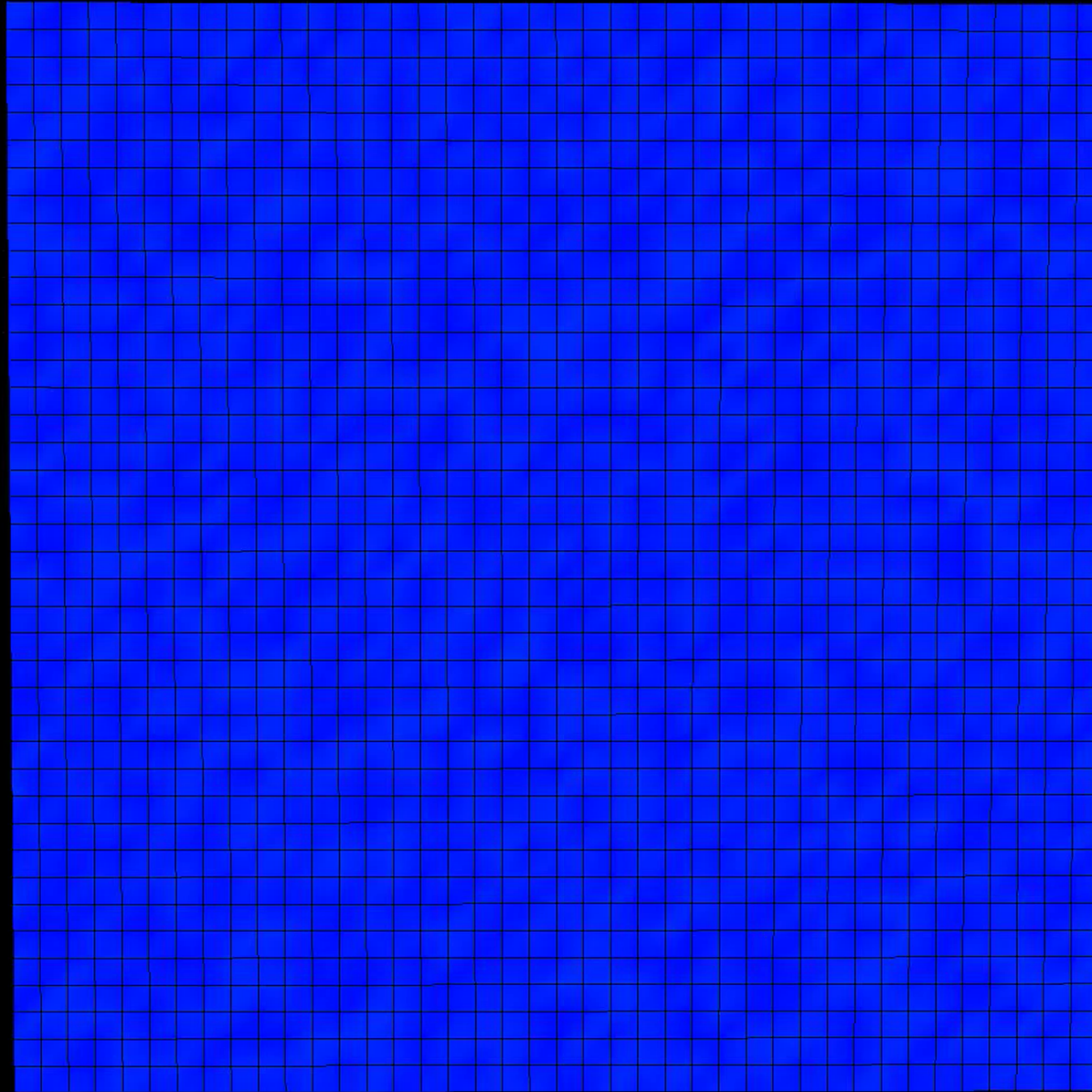


Adaptive Time Step

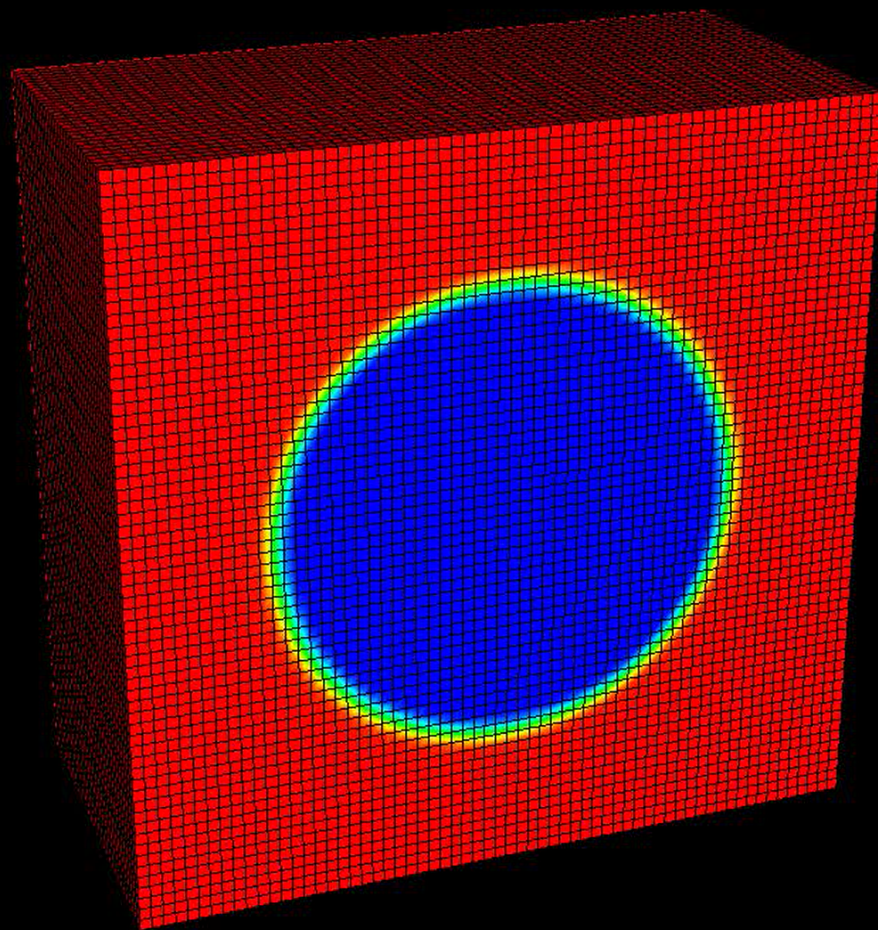
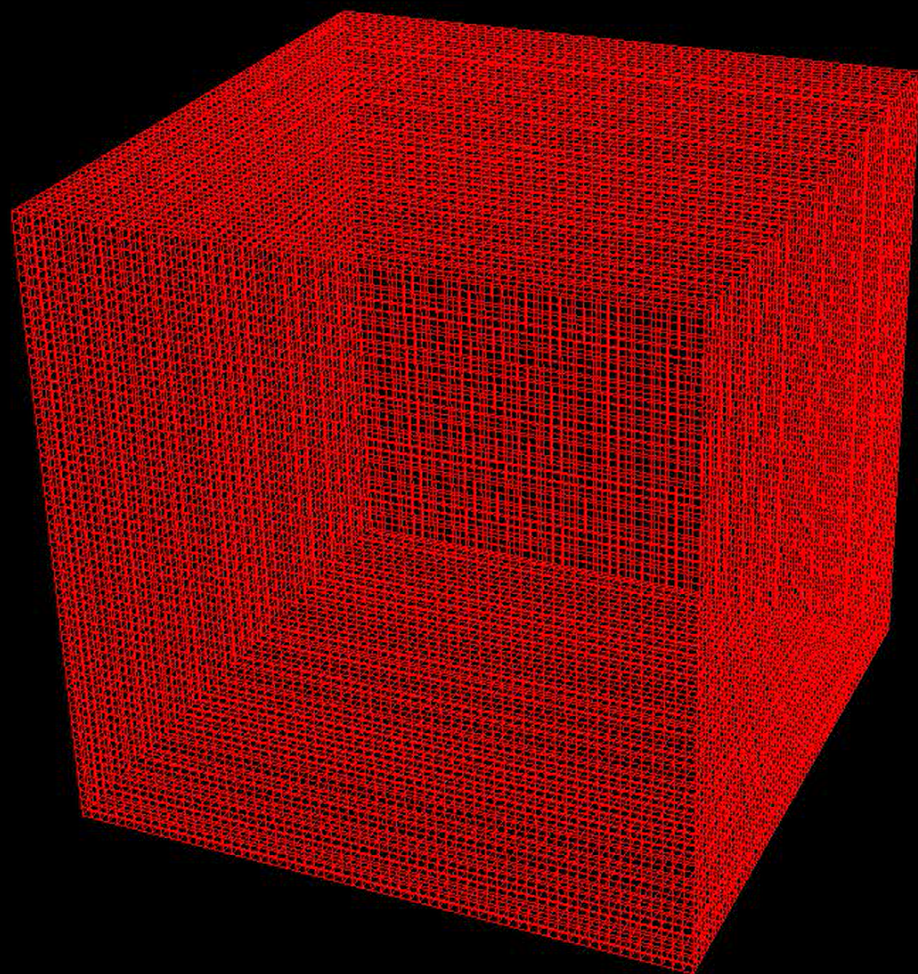
- Solution time step adapts to the current phenomena



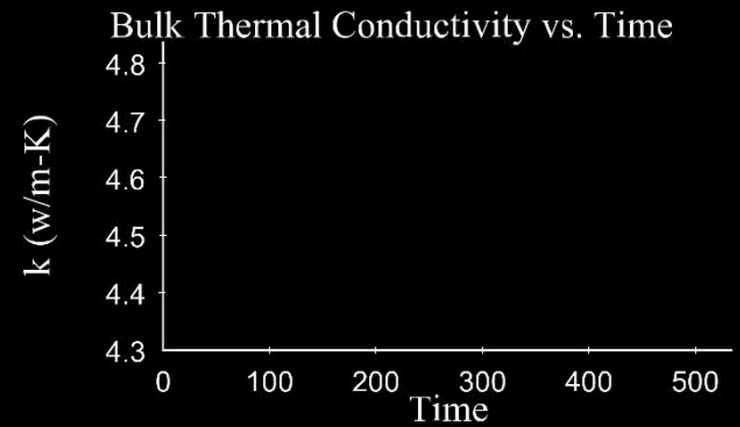
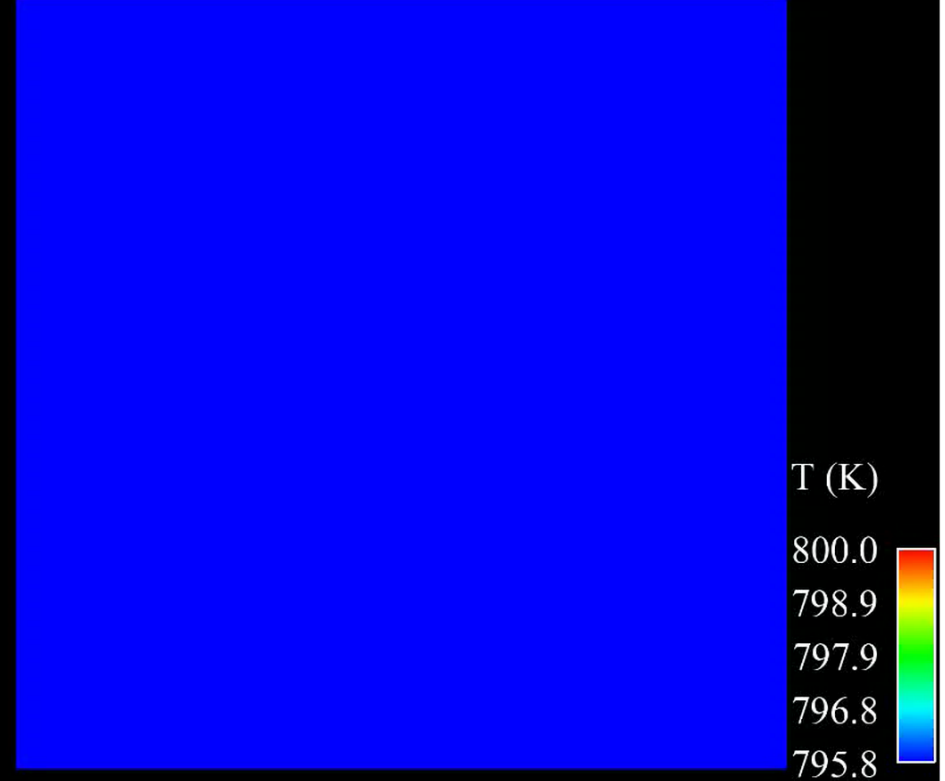
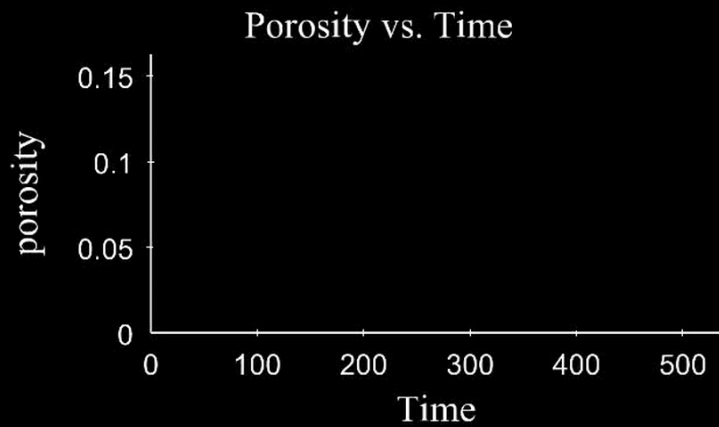
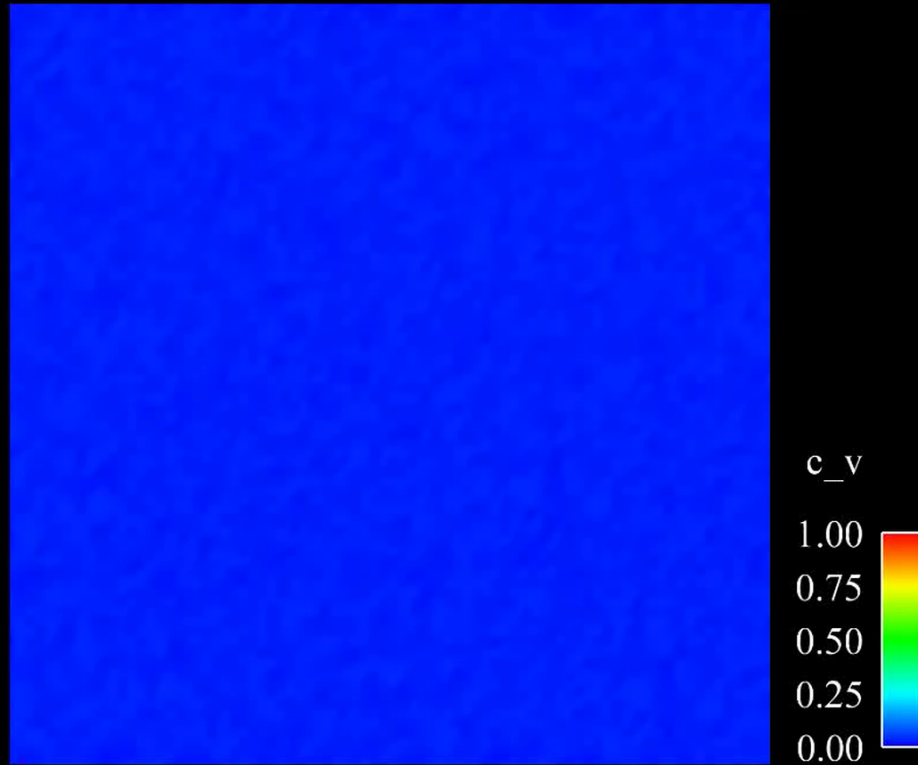
Void Nucleation with Mesh Adaptivity



Shrinking Spherical Grain with Mesh Adaptivity



Effect of Void Nucleation on Thermal Conductivity



Numerical Summary

- Convergence is excellent
- Scalability seems good (we still need to do a full scalability analysis)
- Third-order Hermite elements are expensive. We are currently trying to substitute in an additional variable for the Laplacian interfacial term to allow for linear Lagrange elements.
- Sharp transition in the free energy functional due to the log terms eventually causes problems in the numerical solution

$$f^{solid} = E_v^f c_v + E_i^f c_i + E_g^f c_g + k_B T [c_v \ln(c_v) + c_i \ln(c_i) + c_g \ln(c_g) + (1 - c_v - c_i - c_g) \ln(1 - c_v - c_i - c_g)]$$

